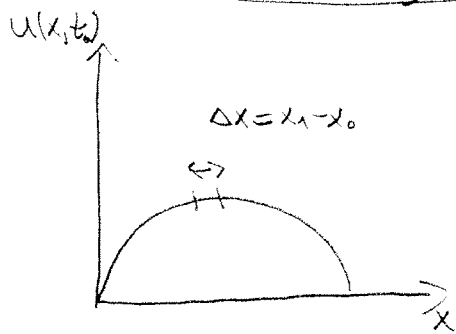


① The violin string (wave equation)

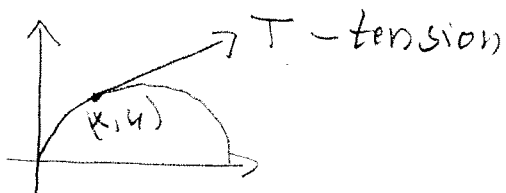


$\rho(x)$ - density
 $m = \rho \Delta x$ - mass
 u_{tt} - acceleration

Newton's law

$$ma = f \text{ - force}$$

The force term



Assume

- * no gravitation (low mass)
- * tension in direction of the tangent \hat{t} (the string does not resist being bent)

Tangent at $x = x_0$ (fix t)

$$u = u(x_0) + u_x(x_0)(x - x_0), \quad \hat{t} = \frac{1}{\sqrt{1+u_x^2}} (1, u_x)$$

$$T = A \hat{t}, \quad A \text{ - amplitude (the tension)}$$

Assume: $T(x) = T; \rho(x) = \rho$

$$\boxed{\text{vertical acceleration } u_{tt}} = \boxed{\text{vertical part of } T}$$

$$f = A \frac{u_x}{\sqrt{1+u_x^2}} \Big|_{x=x_0}^{x=x_1}$$

\Rightarrow

$$\rho u_{tt} = A \frac{1}{\Delta x} \frac{u_x}{\sqrt{1+u_x^2}} \Big|_{x=x_0}^{x=x_1}$$

$$= \frac{\partial}{\partial x} \left(\frac{u_x}{\sqrt{1+u_x^2}} \right) \text{ in the limit } \Delta x = x_1 - x_0 \rightarrow 0$$

The nonlinear wave eq.

$$\rho u_{tt} - A \frac{\partial}{\partial x} \left(\frac{u_x}{\sqrt{1+u_x^2}} \right) = 0$$

(2)

The linear wave eq.

Assume small oscillations: $u_x^2 \ll u_x \ll 1$

$$\Rightarrow \frac{u_x}{\sqrt{1+u_x^2}} = u_x + \mathcal{O}(u_x^3) \approx u_x$$

$$\boxed{u_{tt} - c^2 u_{xx} = 0, \quad c = \frac{A}{\rho}}$$

General solution

Let $\xi = x - ct$, $\eta = x + ct$

\Rightarrow

$$u_x = u_{\xi} \cdot 1 + u_{\eta} \cdot 1 \quad ; \quad u_t = u_{\xi} \cdot (-c) + u_{\eta} \cdot c$$

$$u_{xx} = u_{\xi\xi} + 2u_{\xi\eta} + u_{\eta\eta} \quad ; \quad u_{tt} = c^2 (u_{\xi\xi} - 2u_{\xi\eta} + u_{\eta\eta})$$

$$0 = u_{tt} - c^2 u_{xx} = c^2 (u_{\xi\xi} - 2u_{\xi\eta} + u_{\eta\eta} - u_{\xi\xi} - 2u_{\xi\eta} - u_{\eta\eta}) \\ = -4c^2 u_{\xi\eta}$$

$$\Rightarrow u = \phi(\xi) + \psi(\eta)$$

$$\boxed{u(x,t) = \phi(x-ct) + \psi(x+ct)}$$

right going wave
(outgoing)

left going wave
(incoming)

(3) Time harmonic waves

We want to specify an incoming wave with a fixed frequency and not the data $u(x,0)$ & $u_t(x,0)$

Assume time-harmonic waves:

$$U(x,t) = u(x)e^{-i\omega t} \quad \& \quad U_{tt} - c^2 U_{xx} = 0$$

$$\Rightarrow -\omega^2 u - c^2 u'' = 0 \Leftrightarrow \boxed{u'' + \kappa^2 u = 0, \quad \kappa = \frac{\omega}{c}} \quad \text{Helmholtz eq.}$$

ω - (angular) frequency; c - wave speed; κ - wave number

Plane waves

$$u(x) = Ae^{i\kappa x} + Be^{-i\kappa x}$$

$$u(x+\lambda) = u(x) \quad \text{when} \quad \lambda = \frac{2\pi}{\kappa}$$

λ - wave length

$$U(x,t) = Ae^{i(\kappa x - \omega t)} + Be^{i(\kappa x + \omega t)} \quad \text{- plane waves}$$

constant value if the differential is zero

$$v_p = \frac{dx}{dt} = \frac{\omega}{\kappa} \quad \text{- phase velocity}$$

$e^{i\kappa x}$ - right going or **outgoing** with $v_p = c$

$e^{-i\kappa x}$ - left going or **incoming** with $v_p = -c$

We must eliminate the incoming wave to get an unique solution!

Non-reflecting boundary conditions

Take $u' - i\kappa u = 0$ at $x = x_0$

\Rightarrow eliminate $e^{-i\kappa x}$ from the set $\{e^{i\kappa x}, e^{-i\kappa x}\}$

This do only work in 1D!

(4) Helmholtz eq. in 3D

$$\Delta u + \kappa^2 u = 0, \quad x \in \mathbb{R}^3 \setminus \bar{\Omega} =: \Omega^+$$



No simple non-reflecting boundary condition exist when $\dim > 1$

Sommerfeld's radiation condition

Fundamental solution

$$\Delta G + \kappa^2 G = \delta(x) \Rightarrow G(|x' - x|) = \frac{e^{i\kappa|x' - x|}}{4\pi|x' - x|}$$

Green's theorem \Rightarrow

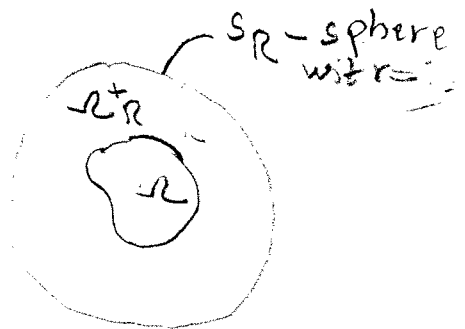
$$u(x) = \int_{\partial\Omega} \left\{ u(x') \frac{\partial}{\partial n(x')} G(x, x') - G(x, x') \frac{\partial}{\partial n(x')} u(x') \right\} dS(x')$$

$F(x, x')$

Truncate the domain:

$$\Omega_R^+ \Rightarrow \Omega^+, \quad R \rightarrow \infty, \quad \partial\Omega_R^+ = \partial\Omega \cup S_R$$

Demand: $\int_{S_R} F(x, x') dS(x') \rightarrow 0, \quad R \rightarrow \infty$



Fix $x \in \Omega^+$: we have $|x' - x| \approx |x'|$ and $\frac{\partial}{\partial n} = \frac{\partial}{\partial R}$ when R large and $x' \in S_R$

$$\Rightarrow \int_{S_R} \frac{1}{R} \left(i\kappa u - \frac{u}{R} - \frac{du}{dR} \right) \frac{e^{i\kappa R}}{4\pi R^2} dS \rightarrow 0, \quad R \rightarrow \infty$$

" \Rightarrow " (a) $Ru < \infty$ for all R

(b) $R \left(i\kappa u - \frac{du}{dR} \right) \rightarrow 0, \quad R \rightarrow \infty$

(i) It can be shown that (b) \rightarrow (a)

(ii) (b) is the Sommerfeld condition