

Sol. Exercise VI

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1)

$$\begin{aligned}
 \text{a)} \quad & \int_{\Omega_R^c} (-\Delta \bar{u} u - \bar{u}^2 |u|^2) ds = \|\nabla u\|_{L^2(\Omega_R^c)}^2 - \bar{u}^2 \|u\|_{L^2(\Omega_R^c)}^2 \\
 & = \int_{\partial \Omega_R^c} \frac{\partial \bar{u}}{\partial n} u ds = \int_{\partial B_R} \frac{\partial \bar{u}}{\partial n} u ds - \int_{\partial \Omega^c} \frac{\partial \bar{u}}{\partial n} u ds
 \end{aligned}$$

b) Multiply by ν and take imaginary parts:

$$\begin{aligned}
 & \text{Im}(\nu) \left(\|\nabla u\|_{L^2(\Omega_R^c)}^2 + |\nu|^2 \|u\|_{L^2(\Omega_R^c)}^2 \right) \\
 & = -\text{Im} \left(\int_{\Gamma} \nu \frac{\partial \bar{u}}{\partial n} u ds \right) + \text{Im} \left(\int_{\partial B_R} \nu \frac{\partial \bar{u}}{\partial n} u ds \right)
 \end{aligned}$$

using (4),

$$\begin{aligned}
 & = -\int_{\Gamma} \text{Im} \left(\nu \frac{\partial \bar{u}}{\partial n} u \right) ds + \frac{1}{2} \int_{\partial B_R} \left| \frac{\partial u}{\partial n} - i\nu u \right|^2 ds \quad (*) \\
 & \quad - \frac{1}{2} \int_{\partial B_R} \left(\left| \frac{\partial u}{\partial n} \right|^2 + |\nu|^2 |u|^2 \right) ds
 \end{aligned}$$

Let $R \rightarrow \infty$. Since $\text{Im}(\nu) > 0$, same field integral goes to zero, and so by the condition on the integral in Γ , $u = 0$ over Ω^c .

c) Now the L.H.S is zero. Then,

$$0 = -\frac{1}{2} \int_{\partial B_R} \left(\left| \frac{\partial u}{\partial n} \right|^2 + |\nu|^2 |u|^2 \right) ds = \int_{\Gamma} \text{Im} \left(\nu \frac{\partial \bar{u}}{\partial n} u \right) ds$$

$\Rightarrow \int_{\Omega^c} |u|^2 ds = 0 \Rightarrow$ Rellich theorem applies $\Rightarrow u = 0 \in \Omega^c$

$$2) \text{ a) } \int_{\partial B_R} u \frac{\partial \bar{u}}{\partial n} ds = \int_{\partial \Omega} u \frac{\partial \bar{u}}{\partial n} ds - u^2 \int_{\Omega_R^c} |u|^2 dx + \int_{\Omega_R^c} |\nabla u|^2 dx \quad (2)$$

Take imaginary parts and limit $R \rightarrow \infty$

$$\int_{\partial B_R} \text{Im} \left(u \frac{\partial \bar{u}}{\partial n} \right) ds = \int_{\Gamma} \text{Im} \left(u \frac{\partial \bar{u}}{\partial n} \right) ds$$

and using (*)

$$\lim_{R \rightarrow \infty} \int_{\partial B_R} \left(\left| \frac{\partial u}{\partial n} \right|^2 + u^2 |u|^2 \right) ds = -2u \text{Im} \int_{\Gamma} u \frac{\partial \bar{u}}{\partial n} ds$$

Since, the l.h.s is composed of positive terms and the r.h.s is bounded. The terms are individually bounded. $\Rightarrow \int_{S_R} |u|^2 ds = o(1)$ as $R \rightarrow \infty$.

b)

$$|I_1| \leq \left(\int_{S_R} |u|^2 ds \right)^{1/2} \cdot \left(\int_{S_R} \underbrace{\left| \frac{\partial \phi}{\partial n} - iu\phi \right|^2}_{o\left(\frac{1}{R^2}\right)} ds(y) \right)^{1/2} \rightarrow 0 \text{ as } R \rightarrow \infty.$$

$$|I_2| \leq \left(\int_{S_R} \underbrace{|\phi(x,y)|^2}_{o\left(\frac{1}{R}\right)} ds(y) \right)^{1/2} \cdot \left(\int_{S_R} \underbrace{\left| \frac{\partial u}{\partial n} - iuu \right|^2}_{o\left(\frac{1}{R^2}\right)} ds(y) \right)^{1/2} \rightarrow 0 \text{ as } R \rightarrow \infty.$$

c) From Green's second inequality,

$$\int_{\Omega_R^c} (-\Delta u - u^2 u) \phi(x,y) + (\Delta \phi + u^2 \phi) u =$$

$$= \int_{\partial B_R^c} \underbrace{\left(u \frac{\partial b}{\partial u} - \frac{\partial u}{\partial u} b \right)}_{F(u,b)} ds = \int_{\partial B_R} + \int_{\partial \Omega} \quad (3)$$

From the representation formula in a bounded domain:

$$u(x) = \int_{\partial B_R} F(u,b) + \int_{\Omega} F(u,b)$$

Part

$$\int_{\partial B_R} F(u,b) = I_1 - I_2 \rightarrow 0 \text{ as } R \rightarrow \infty.$$

3] a) Taylor $u = 1, 0 < x < 1$

$$\frac{1}{\sqrt{1-xe^{\pm\theta^2}}} = \frac{1}{\sqrt{1-x}} = 1 + \left(\frac{1}{2} \frac{(-1)}{(\sqrt{1-x})^3} \right) \Big|_{x=0} x$$

$$= 1 + \frac{1}{2} x$$

$u=2$

$$\frac{d}{dx} \left(\frac{1}{2} (1-x)^{-3/2} \right) = \frac{1}{2} \left(-\frac{3}{2} \right) (1-x)^{-5/2} \cdot (-1)$$

$$\Rightarrow \frac{d^2}{dx^2} \left(\frac{1}{2} (1-x)^{-3/2} \right) = \frac{1}{2} \cdot \frac{3}{2} (1-x)^{-5/2} \Rightarrow \frac{1}{\sqrt{1-x}} = 1 + \frac{1}{2} x + \frac{3}{4} x^2$$

Induction:

$$\frac{d^u}{dx^u} (1-x)^{-1/2} = \frac{1 \cdot 3 \cdots (2u-1)}{2 \cdot 4 \cdots 2u} (1-x)^{-\frac{(2u+1)}{2}}$$

$$\frac{d}{dx} \left(\frac{d^u}{dx^u} (1-x)^{-1/2} \right) = \frac{1 \cdots (2u-1)}{2 \cdots 2u} (1-x)^{-\frac{(2u+1)}{2}} \cdot (-1)$$

$$= \frac{d^{u+1}}{dx^{u+1}} (1-x)^{-1/2} \cdot \frac{-(2u+3)}{2}$$

Ratio test:

$$\lim_{n \rightarrow \infty} \left| \frac{a_{n+1}}{a_n} \right| = \frac{2n+1}{2(n+1)} \cdot r < 1$$

\Rightarrow converges absolutely

M-test:

$$\sum_{n=1}^{\infty} \frac{1 \cdot 3 \cdots (2n-1)}{2 \cdot 4 \cdots 2n} r^n e^{t \cos \theta} \leq \sum_{n=1}^{\infty} \left| \frac{1 \cdot 3 \cdots (2n-1)}{2 \cdot 4 \cdots 2n} \right| r^n$$

\rightarrow uniformly convergent

$$\leq \sum_{n=1}^{\infty} \left| \frac{1 \cdot 2 \cdots n}{1 \cdot 2 \cdots n} \right| r_0^n$$

b) Notice that

$$r_0 \in (0, 1)$$

$$\frac{1}{\sqrt{1-r e^{i\theta}}} \cdot \frac{1}{\sqrt{1-r e^{-i\theta}}} = \frac{1}{\sqrt{1-2r \cos \theta + r^2}} \quad z = \cos \theta$$

Since the series are convergent (uniformly and absolutely) take $\theta = 0$. Then, the geometric series is a support for (9)

$$\rightarrow |P_n(t)| \geq 1 \quad -1 \leq t \leq 1$$

c) Differentiating (9) yields w.r.t. t

$$-\frac{1}{2} \frac{(-2t+2r)}{(1-2t+r+r^2)^{3/2}} = \sum_{n=0}^{\infty} n P_n(t) r^{n-1} / (1-2t+r+r^2)$$

$$\frac{t-r}{\sqrt{1-2t+r+r^2}} = \sum_{n=0}^{\infty} n P_n(t) r^{n-1} (1-2t+r+r^2)$$

$$\sum_{n=0}^{\infty} P_n(t) r^n (t-r) = \sum_{n=0}^{\infty} n P_n(t) r^{n-1} (1-2t+r+r^2)$$

Coefficient: r^u

$$P_n t^2 - P_{n+1} = n P_n (-2t) + n P_{n-1} + n P_{n+1}$$

$$\Rightarrow 0 = (n+1) P_{n+1}(t) - (2n+1)t P_n(t) + n P_{n-1}$$

d) $\int_{-1}^1 \frac{dt}{1-2t+r^2} = \int_{-1}^1 \left(\sum_{n=0}^{\infty} P_n(t) r^n \right)^2 dt$ $n=0, 1, \dots$

By orthogonality of P_n (spherical harmonics)

$$\int_{-1}^1 g^2 dt = \sum_{n=0}^{\infty} \int_{-1}^1 [P_n(t)]^2 dt r^{2n}$$

uniform convergence allows the exchange of integration/sum
On the other hand,

$$\int_{-1}^1 \frac{dt}{1-2t+r^2} = \frac{1}{r} \ln \frac{1+r}{1-r} = \sum_{n=0}^{\infty} \frac{2}{2n+1} r^{2n}$$

Then $\int_{-1}^1 [P_n(t)]^2 dt = \frac{2}{2n+1}$

4) a) Each harmonic polynomial is mapped again into homogeneous harmonic polynomials by \mathcal{Q} orthogonal $\Rightarrow (k')$
Now, since \mathcal{Q} is orthogonal and Y_n^m are orthogonal

$$\int_{\mathbb{S}^2} Y_n^m(\mathcal{Q}\hat{x}) \overline{Y_n^{m'}}(\mathcal{Q}\hat{x}) ds = \int_{\mathbb{S}^2} Y_n^m(\hat{k}) \overline{Y_n^{m'}}(\hat{k}) ds = \delta_{mm'}$$

Here, in $Y_n^m(\theta, \phi) = \sum_{k=-n}^n a_{mk} Y_n^k(\hat{x})$ $m = -n, \dots, n$

$A = (a_{mk})$ is orthogonal and

$$\begin{aligned} Y(\hat{x}, \hat{y}) &= \sum_{m=-n}^n \left(\sum_{k=-n}^n a_{mk} Y_n^k(\hat{x}) \right) \overline{\left(\sum_{l=-n}^n a_{ml} Y_n^l(\hat{y}) \right)} \\ &= \sum_{m=-n}^n Y_n^m(\hat{x}) \overline{Y_n^m(\hat{y})} = Y(\hat{x}, \hat{y}) \end{aligned}$$

$$\Rightarrow Y(\hat{x}, \hat{y}) = f(\cos\theta)$$

$$b) \sum_{m=-n}^n Y_n^m(\hat{x}) \overline{Y_n^m(\hat{y})} = a_n P_n(\cos\theta)$$

$$\text{Set } \hat{x} = \hat{y} \Rightarrow \theta = 0, \cos\theta = 1$$

$$\Rightarrow \sum_{m=-n}^n |Y_n^m(\hat{x})|^2 = a_n P_n(1) = a_n / \int_{\mathbb{S}^2} d\Omega$$

$$\Rightarrow 4\pi a_n = \sum_{m=-n}^n \int_{\mathbb{S}^2} |Y_n^m(\hat{x})|^2 d\Omega$$

$$\Rightarrow a_n = \frac{2n+1}{4\pi} \quad \Bigg| \quad = 1 \quad \text{by definition.}$$